

X-RAY SPECTRAL STUDIES OF THE ELECTRONIC STRUCTURE OF URANYL FLUORITE UO_2F_2

by

Igor O. UTKIN¹, Yury A. TETERIN¹, Vladimir A. TEREHOV², Mikhail V. RYZHKOV³,
Anton Yu. TETERIN¹, and Labud VUKCHEVICH⁴

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This work interpreted the fine X-ray photoelectron spectral structure of the low binding energy electrons (0-40 eV) and X-ray $\text{O}_{4,5}(\text{U})$ emission spectral structure from UO_2F_2 taking into account the relativistic X_α discrete variation ($\text{RX}_\alpha\text{-DV}$) calculation for the $[(\text{UO}_2)\text{F}_6]^{4-}(\text{D}_{6h})$ cluster reflecting an uranium close environment in UO_2F_2 . The U5f electrons were shown to participate directly in the chemical bond formation. The U6p electrons were shown to participate not only in formation of the inner valence molecular orbitals, but also in formation of the outer valence molecular orbitals. The inner valence molecular orbitals sequence order in the binding energy range 12-40 eV was established. It is important for development of the technique of interatomic distance determination in the axial direction and equatorial plane of uranyl compounds on the X-ray photoelectron spectral basis.

Key words: X-ray photoelectron spectroscopy, outer and inner valence molecular orbitals

INTRODUCTION

The low binding energy (0-50 eV) X-ray photoelectron spectroscopy (XPS) from uranyl fluorite (UO_2F_2) exhibits a complex fine structure attributed to the outer valence molecular orbitals (OVMO) and inner valence molecular orbitals (IVMO) electrons [1-3]. This spectrum reflects the band structure and is observed as several [eV] wide bands. In UO_2F_2 these IVMOs are built from the U6p and O(F)2s shells. A strong U6p-O(F)2s atomic orbital (AO) overlap for the $[(\text{UO}_2)\text{F}_6]^{6-}(\text{D}_{6h})$ cluster in the non-relativistic approximation was theoretically shown in [4, 5] and for the first time experimentally confirmed on the basis of the outer-core shells binding energy differ-

ences for uranyl fluorite in [1, 6]. It allowed a qualitative identification of the high resolution X-ray $\text{O}_{4,5}(\text{U})$ emission spectral structure from UO_2F_2 [4, 7], which appeared to be an extra experimental corroboration for the IVMO formation in UO_2F_2 .

According to earlier suggestions, the An5f electrons are supposed to get promoted to, for example, the An6d atomic orbitals before the chemical bond formation. However, the theoretical calculations show that the An5f atomic shells can participate directly in the formation of molecular orbitals (MO) in actinide compounds [1]. This principally novel fact needs an experimental confirmation so in this work we measured high resolution X-ray $\text{O}_{4,5}(\text{U})$ emission spectral structures. However, in the absence of relativistic calculations of the UO_2F_2 electronic structure a correct interpretation of the XPS can not be achieved. In this work we interpreted the fine low binding energy XPS and X-ray emission spectral (XES) structures from UO_2F_2 taking into account the relativistic X_α discrete variation method ($\text{RX}_\alpha\text{-DVM}$) calculation of the electronic structure for the $[(\text{UO}_2)\text{F}_6]^{4-}(\text{D}_{6h})$ cluster reflecting an uranium close environment in UO_2F_2 .

EXPERIMENTAL

The XPS method is based on the photoeffect phenomenon in which XPS spectrum shows a dependence

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Authors' addresses:

¹ Russian Research Center "Kurchatov Institute"
1, Kurchatov square, Moscow 123182, Russia

² Voronezh State University, Voronezh, Russia

³ Institute of Solid-State Chemistry of Ural Dept. of RAS,
Ekaterinburg, Russia

⁴ Faculty of Natural Sciences and Mathematics,
University of Montenegro
P. O. Box 211, 81000 Podgorica, Serbia and Montenegro

E-mail address of corresponding author:
utkin@ignph.kiae.ru (I. O. Utkin)

of the number of photoelectrons on the binding energy. XPS spectra of UO_2F_2 were measured with an electrostatic spectrometer HP 5950A Hewlett-Packard using monochromatic $\text{AlK}_{\alpha,1,2}$ ($h\nu = 1486.6$ eV) radiation in a vacuum of $1.3 \cdot 10^{-7}$ Pa at room temperature. The device resolution measured as full width on the half-maximum (FWHM) of the $\text{Au}4f_{7/2}$ line on the standard rectangular golden plate was 0.8 eV. The binding energies E_b [eV] were measured relative to the binding energy of the Cls electrons from hydrocarbons adsorbed on the sample surface accepted to be equal to 285.0 eV. The FWHM were measured relative to the width of the Cls line of hydrocarbons accepted to be equal to 1.3 eV. The error in determination of electron binding energies and the line widths did not exceed 0.1 eV and that of the relative line intensities was less than 10%.

UO_2F_2 sample was prepared from the fine-dispersed powder ground in the agate mortar as a thick dense layer with a flat surface pressed in indium on the metallic substrate. Elastic scattering related background in the XPS was subtracted as described by Shirley [8].

X-ray $\text{O}_{4,5}(\text{U})$ emission spectra (XES) of UO_2F_2 reflecting the U6p and U5f states in the valence band were measured by a primary method with a spectrometer RSM-500 with an energy resolution of 0.3 eV. The spectra were recorded during 60 minutes [4]. The change of energy of incident electrons allowed varying an effective depth of analysis from 15 up to 50 nm. The area of an exited electronic beam was 5×5 mm². The samples were ground in the agate mortar and then as a powder were pressed in a silver plate, which was fastened on the anode of the X-ray tube with the special clamps. For the spectra registration the secondary electronic multiplier of an open type was used. As the photocathode a CsI film was applied. The spectra were recorded three times to avoid the influence of the sample composition changes. For calibration of energy position of spectrum peculiarities Zr (75.55 eV) and Nb (85.85 eV) lines were taken in the second order of reflection M.

The $[(\text{UO}_2)\text{F}_6]^{6-}(\text{D}_{6h})$ cluster [9] reflecting uranium close environment in UO_2F_2 is a UO_2^{2+} ion surrounded in the equatorial plane by the six fluorines (ligands). The interatomic distances used in the calculations were $R_{\text{U-O}} = 1.74 \cdot 10^{-10}$ m and $R_{\text{U-L(F)}} = 2.43 \cdot 10^{-10}$ m. The calculations for this cluster were done in relativistic X_α Discrete Variation (RX $_\alpha$ -DV) approximation [10] based on the Dirac-Slater equation for the 4 component spinors being a linear combination of atomic orbitals from the solution of the relativistic task for the isolated atoms. The extended basis of numerical atomic orbitals from the solution of Dirac-Slater equation for the isolated atoms beside the filled included the vacant U7p states. The direct AO decomposition basis were symmetrized linear AO combinations converting by the irreducible representations of the binary groups D_{6h} . The numerical integration during the calculation of secular equation matrix elements was done by the number

of 4500 points spread in the cluster space. It provided the convergence of MO energies of not worse than 0.1 eV. The local exchange-correlation potential was taken as X_α with α equal to the mean atomic values. Since the clusters were the fragments of the crystal, the ligand AO population renormalization during the self-consistent was done. It allowed an effective account of the stoichiometry and charge redistribution between the ligands and surrounding crystal [10]. For the correct evaluation of the relativistic effects in UO_2F_2 the non-relativistic X_α -DV calculations (NRX $_\alpha$ -DV) for the same crystalline parameters were done [4, 10].

RESULTS AND DISCUSSION

The low binding energy (0-40 eV) XPS from UO_2F_2 can be conditionally subdivided into the two ranges (fig. 1). The first one 0-12 eV shows the structure attributed to the OVMOs built mostly from the incompletely filled outer U5f,6d,7s and

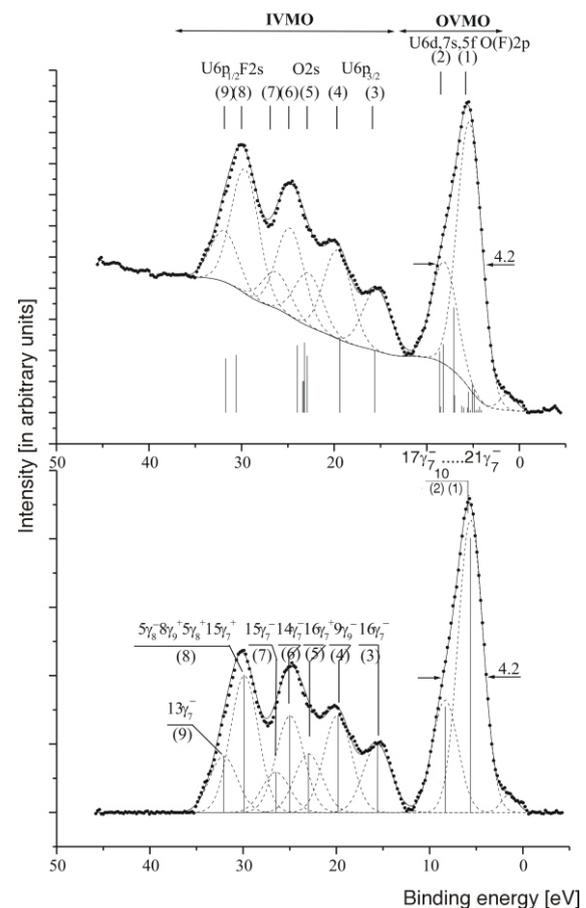


Figure 1. XPS from UO_2F_2 (a). The corresponding theoretical spectrum is given below as vertical bars. The same spectrum with subtracted background (b). Vertical bars show the spectrum on the basis of the theoretical and experimental data (see tab. 2). The dashed line shows the shape of the subtracted background (a) and subdivision into the separate components. Spectral intensity is in arbitrary units normalized by 100%

Table 1. Relativistic X_α -DV calculation results of the electronic structure for the $\text{UO}_2\text{F}_6^{4-}(\text{D}_{6h})$ cluster reflecting uranium close environment in UO_2F_2 at $R_{\text{U-O}} = 1.74 \cdot 10^{-10}$ m and $R_{\text{U-L}} = 2.43 \cdot 10^{-10}$ m

MO	-E [eV]	MO composition																		
		6s $i^{(a)}$ 1.14	6p _{1/2} 0.89	6p _{3/2} 1.29	6d _{3/2} 0.61	6d _{5/2} 0.55	7s 0.12	5f _{5/2} 3.67	5f _{7/2} 3.48	7p _{1/2} 0.07	7p _{3/2} 0.10	O2s 0.96	O2p 0.07	F2s 1.44	F2p 0.13					
OVMO	24 γ^+	-13.0				0.33	0.52	0.14												
	γ^+	-6.0				0.56	0.39	0.04												
	25 γ^-	-5.7										1.00								
	14 ρ^+	-5.7																		
	22 γ^+	-5.3						1.00												
	24 γ^-	-4.5																		
	23 γ^-	-3.9			0.33															
	16 ρ^-	-2.9								0.16	0.43									
	13 ρ^+	-1.0					0.77	0.18												
	9 δ^+	-0.9						0.96												
	15 ρ^-	0.1																		
	22 γ^-	0.3								0.10	0.72									
	11 δ^-	1.1								0.54	0.30									
	10 δ^-	1.4								0.04	0.94									
	9 δ^-	2.1								0.91	0.05									
	14 ρ^-	2.2								0.80	0.14									
	$I_i^{(a)}$				0.33	1.66	3.05	1.18	2.55	3.56	1.00	2.05		0.37				0.22		
	IVMO	21 $\gamma^{+(b)}$	2.9																	1.00
		γ^-	3.1			0.04														0.95
		21 γ^-	3.3																	1.00
8 δ^+		3.5																	1.00	
8 δ^-		3.7																	0.97	
12 ρ^-		3.7																	0.95	
20 γ^-		3.9			0.08														0.81	
7 δ^+		3.9																	1.00	
12 ρ^+		3.9																	0.96	
20 γ^+		3.9																	0.92	
11 ρ^+		4.1																	1.00	
11 ρ^-		4.3																	1.00	
7 δ^-		4.3																	0.93	
19 γ^-		4.4			0.01															0.97
19 γ^+		4.8	0.03					0.02	0.01											0.91
6 δ^-		5.0																		0.98
6 δ^+		5.8								0.18	0.02									0.82
10 ρ^+		5.8								0.19										0.81
18 γ^-		5.9			0.17															0.19
18 γ^+		7.0	0.07			0.02	0.04			0.12	0.13				0.05					0.04
10 ρ^-	7.0			0.06															0.71	
9 ρ^+	7.3								0.07	0.03									0.90	
17 γ^-	7.4			0.03															0.75	
17 γ^+	7.4								0.18										0.87	
$\Sigma I_i^{(a)}$		0.10	0.04	0.35	0.02	0.55	0.02	0.40	0.63			0.05	4.47						17.3	
IVMO	16 γ^-	14.4		0.03	0.47								0.30	0.20						
	9 ρ^-	18.2			0.87								0.05	0.06					0.01	
	16 γ^+	21.7	0.10			0.01	0.01						0.85	0.03						
	15 γ^-	22.0		0.12	0.04								0.07	0.76						
	5 δ^-	22.0												1.00						
	8 ρ^+	22.1												1.00						
	5 δ^+	22.2												1.00						
	15 γ^+	22.8												0.96						
	8 ρ^-	22.8			0.07									0.93						
14 γ^-	24.4		0.26	0.24								0.24	0.04							
13 γ^-	30.5		0.56	0.04								0.29	0.10							
14 γ^+	46.4	0.80										0.15	0.04							
$\Sigma I_i^{(a)}$		1.00	0.97	1.73	0.01	0.01						1.93	0.44	5.97	0.01					

^(a) Total electronic densities per one electron^(b) Upper filled molecular orbital

O(F)2p AOs (tab. 1). The second one 12-40 eV shows the IVMO related fine structure. These IVMOs are built mostly from the completely filled inner valence U6p and O(F)2s AOs. The fact that the IVMO XPS parameters correlate with uranium close environment structure in compounds encour-

aged this subdivision [11]. The OVMO XPS structure has three typical features and can be subdivided into the three components. Since the low binding energy component at $E_b = 1.5$ eV is apparently due to the admixture of the U^{4+} ions, it will not be considered here. The IVMO spectral range exhibits

pronounced peaks and can be subdivided into the seven components (fig. 1). Despite a certain formalism, this subdivision allows qualitative and quantitative comparisons between the XPS, XES, and RX_{α} -DV calculation results for the $[(UO_2)F_6]^{6-}$ (D_{6h}) cluster.

The relativistic calculation results for the ground state of this cluster is given in tab. 1. Since after the electron photoemission the molecule transits into an excited state with a hole on a certain level, the binding energies for the transition states have to be used [12] for a more strict comparison between the theoretical and experimental data. However, in a rough approximation one can suggest that for the valence region the binding energies for the transition state differ from those for the ground state by a constant shift toward the higher absolute energy. Therefore, the present work used the theoretical binding energies increased by 1.2 eV for comparison with the corresponding experimental values (tab. 2). Taking into account the MO compositions (tab. 1) and photoionization cross-sections [13], the theoretical spectral intensities for some energy ranges were determined (tab. 2, fig. 1). Comparing the experimental XPS and theoretical data one has to take into account that the XPS from UO_2F_2 reflects the band structure and consists of bands widened due to the solid-state effects. Despite the approximation imperfections, one can see a good qualitative agreement between the theoretical and experimental data. Indeed, the corresponding theoretical and experimental FWHMs ($\Gamma_{\text{theor}} = 4.5$ eV and $\Gamma_{\text{exp}} = 4.2$ eV) and relative intensities of the outer and inner valence bands ($I_{OVMO}^{\text{theor}} = 42.0\%$, $I_{IVMO}^{\text{theor}} = 58.0\%$ and $I_{OVMO}^{\text{exp}} = 40.6\%$, $I_{IVMO}^{\text{exp}} = 59.4\%$) are comparable. A satisfactory agreement between the theoretical and experimental binding energies for some electronic orbitals, in particular the $16 \tau^-$, $9 \rho^-$, and $16 \tau^+$ IVMO, was reached (tab. 2). The worst agreement was reached for the middle ($15 \tau^-$ - $14 \tau^-$) IVMO region. Previously the IVMO XPS structure has been interpreted on the basis of the non-relativistic X_{α} calculations [4, 5]. It allowed a qualitative identification of the fine spectral structure in the binding energy range 0-23 eV. It was possible because inclusion of relativistic effects in general gives a strong separation of the $U6p_{1/2}$ component, but the main features in the binding energy range 0-23 eV remained invariable. The relativistic calculation allowed a qualitative and quantitative interpretation of the XPS fine structure in the whole range 0-40 eV.

Despite the fact that the outer valence band intensity for uranium compounds is mostly due to the $U5f, 6d$ and $O(F)2p$ AOs, for uranium oxides UO_2 , UO_3 , and fluorite UF_4 it has been shown on the basis of the conversion electron spectra that the inner valence $U6p$ electrons participate effectively in the OVMO formation [14-16]. Using this and the data of tab. 2 in mind one can suggest that the $U6p$ elec-

trons participate in the OVMO formation in UO_2F_2 . Also, experimental evidence for the fact that the $U5f$ electrons can participate in the chemical bond without losing their f-nature was practically established. Indeed, the experimental intensity ratio OVMO/IVMO is 0.68, which is comparable with the corresponding theoretical value 0.72 (tab. 2). Since the $U5f$ AO practically does not participate in the IVMO formation, the difference can be explained by the decrease of the photoionization cross-section of the $U5f$ electrons participating in the chemical bond. It contradicts the traditional in chemistry by suggesting the excitation of the $U5f$ electrons to the $U6d$ shell before the chemical bond formation. Therefore, one can suggest that the relativistic calculation results reflect the partial $U5f$ electronic density rather well (tab. 2). These data show that more than two $U5f$ electrons participate in the chemical bond formation, and the vacant $U5f$ states are located within 6 eV near the absorption edge (tab. 1).

In the IVMO XPS range the good agreement was reached only for the $16 \tau^-$, $9 \rho^-$, $16 \tau^+$, and $13 \tau^-$ orbitals characterizing the spectral width. Taking into account the comparability of the theoretical data and experimental total relative intensity, one can suggest that the calculated energies of the $15 \tau^-$ - $14 \tau^-$ IVMO differ significantly from the corresponding experimental data (tab. 2).

Taking into account the non-relativistic calculations for the $[(UO_2)F_6]^{4-}$ (D_{6h}) cluster [7] and experimental data on the core-valence levels binding energy differences for uranium [17] and UO_2F_2 [1-4], in the molecular orbitals as linear combinations of atomic orbitals approximative MO schematic diagram can be built (fig. 2). This diagram enables the understanding of the real XPS structure of UO_2F_2 . In this approximation one can separate formally the $16\gamma_7^-(3)$, $14\gamma_7^-(6)$, and $13\gamma_7^-(9)$ IVMO group characterizing the interatomic distance U-O in uranyl group UO_2^{2+} , $9\rho_9^-(4)$, $15\gamma_7^-(7)$, and $8\rho_9^-(8)$ IVMO characterizing the interatomic distance U-L in the equatorial plane, as well as the quasiautomatic $16\gamma_7^+(5)$, $5\gamma_8^-(8)$, $8\rho_9^+(8)$, $5\gamma_8^+(8)$, and $15\gamma_7^+(8)$ IVMOs attributed generally to the $O(F)2s-U6p_{1/2}$ electrons. It has to be noted that this IVMO subdivision is well grounded [1, 6]. The experimental data show that the binding energies of quasiautomatic $5\gamma_8^-(8)$ - $15\gamma_7^+(8)$ IVMOs have to be close to 29.9 eV, which differs significantly from results of theoretical calculation. Indeed, the F1s and F2s binding energies for UO_2F_2 are 685.3 and 29.9 eV respectively and $\Delta E_i(F1s-F2s) = 655.4$ eV [1, 6]. In support of this one can note the similar ΔE_i values of 655.4 and 655.2 eV for MnF_2 and MnF_3 respectively [18] where the F2s AOs participate weakly in the MO formation. It allowed a more exact determination of the IVMO sequence order (compare tab. 1 and tab. 2, fig. 2). Since the binding energies of the quasiautomatic F2s related orbitals increase by

Table 2. XPS parameters of the UO₂F₂ and UO₂F₆⁴⁻ (D_{6h}) cluster at $R_{U-O} = 1.74 \cdot 10^{-10}$ m and $R_{U-O} = 2.43 \cdot 10^{-10}$ m ($R_{X\alpha}$ -DVM) and the U6p and U5f electronic state densities $\rho_i(e^-)$

MO		-E ^(c) [eV]	X-ray photoelectron spectra			U6p and U5f electronic state densities $\rho_i(e^-)$				
			Energy ^(d) [eV]	Intensity [%]		5f _{5/2}	5f _{7/2}	6p _{3/2}	6p _{1/2}	
			Experiment	Theory	Experiment					
OVMO	21 γ^+ ^(a)	4.1		0.3						
	13 γ^-	4.3		0.7						0.08
	21 γ^-	4.5		0.3						
	8 8^+	4.7		0.3			0.06			
	8 8^-	4.9		1.1			0.10			
	12 γ^-	4.9		1.6			0.12			
	20 γ^-	5.1		2.6						0.16
	7 8^+	5.1		0.3						
	12 γ^+	5.1		0.3						
	20 γ^+	5.1		0.3						
	11 γ^+	5.3		0.3						
	11 γ^-	5.5		0.3				0.14		
	7 8^-	5.5		2.1						
	19 γ^-	5.6		0.6					0.02	
	19 γ^+	6.0		0.6			0.02			
	6 8^-	6.2		0.8						
	6 8^+	7.0		1.0			0.04			
	10 γ^+	7.0		1.0				0.52		
	18 γ^-	7.1	5.6 (3.1)	12.0		29.3				0.34
18 γ^+	8.2		1.6			0.24	0.26			
10 γ^-	8.2	8.3 (3.2)	6.2		11.3				0.12	
9 γ^+	8.5		0.7			0.14	0.06			
17 γ^-	8.6		6.3					0.06		
17 γ^+	8.6		0.7			0.36				
	$I_i^{(b)}$			42.0	40.6	0.80	1.26	0.08	0.70	
IVMO	16 γ^-	15.6	15.6(3.5)	7.1	7.3			0.06	0.94	
	9 γ^-	19.4	19.8(3.5)	8.7	10.4				1.74	
	16 γ^+	22.9	23.0(3.5)	6.5	6.3					
	15 γ^-	23.2	26.5(2.5)	4.4	4.3			0.24	0.08	
	5 8^-	23.2		3.6						
	8 γ^+	23.3		3.6						
	5 8^+	23.4	29.9(3.5)	3.6	14.7					
	15 γ^+	24.0		3.7						
	8 γ^-	24.0		4.0					0.14	
	14 γ^-	30.6	25.0(3.5)	6.6	10.4			0.52	0.48	
	13 γ^-	31.7	32.1(3.5)	6.2	6.0			1.12	0.08	
	$I_i^{(b)}$	47.6	49.3(6.0)	58.0	59.4			1.94	3.46	
	14 γ^+			~8.0						

(a) Upper filled MO (2 electrons), filling number for all orbitals is 2

(b) Total U6p electronic state densities and peak intensities

(c) Binding energies shifted by 1.2 eV toward the higher values

(d) FWHMs are given in parentheses

about 6 eV, the F2s contribution to, for example, the $9\gamma_9^-(4)$ IVMO must decrease, and to the $13\gamma_7^-(9)$ IVMO must increase. The contribution of the U6p_{3/2} AO to the $8\gamma_9^-(8)$ IVMO must decrease, which agrees with the experimental data. In this case the separation of the $8\gamma_9^-(8)$ IVMO from the general band of quasiatomic orbitals (8) is not possible (tab. 2). Taking into account that $\Delta E_U = 360.6$ eV, $\Delta E_1 = 363.6$ eV, one can find $\Delta_2 = 3.0$ eV and $\Delta_1 = 4.2$ eV [1, 6]. The perturbation is $\Delta_3 = 2.6$ eV. The peak width of the quasiatomic orbitals $16\gamma_7^+(5)$ and $5\gamma_8^-(8)$ (8) is comparable to the widths of other IVMOs. The observed narrowing of the $15\gamma_7^-(7)$ IVMO can be explained by the partial loss of the antibonding nature due to the mixture of the (7%) of O2s AO (tab. 2, see also [11]).

For comparative quantitative analysis of the experimental and theoretical intensities the considered

spectra were decomposed having in mind data in fig. 2. The diagram of fig. 2 was built on the basis of the experimental binding energies and theoretical intensities. The identification of the XPS structure is shown in tab. 2 and fig. 2. These data show that the experimental and theoretical binding energies practically coincide. However, the experimental IVMO intensities often differ significantly from the corresponding theoretical values. The best agreement was reached for the $16\gamma_7^-(3)$, $16\gamma_7^+(5)$, and $13\gamma_7^-(9)$ IVMOs. The theoretical data show that 0.78 U6p e^- electrons participate in the OVMO formation (tab. 2), where the U6p_{3/2} electrons take the main part. The U6p distribution in OVMO and IVMO reflects the distribution of the U6p partial electronic density in the valence band of UO₂F₂. Table 2 also shows that more than two U5f electrons partici-

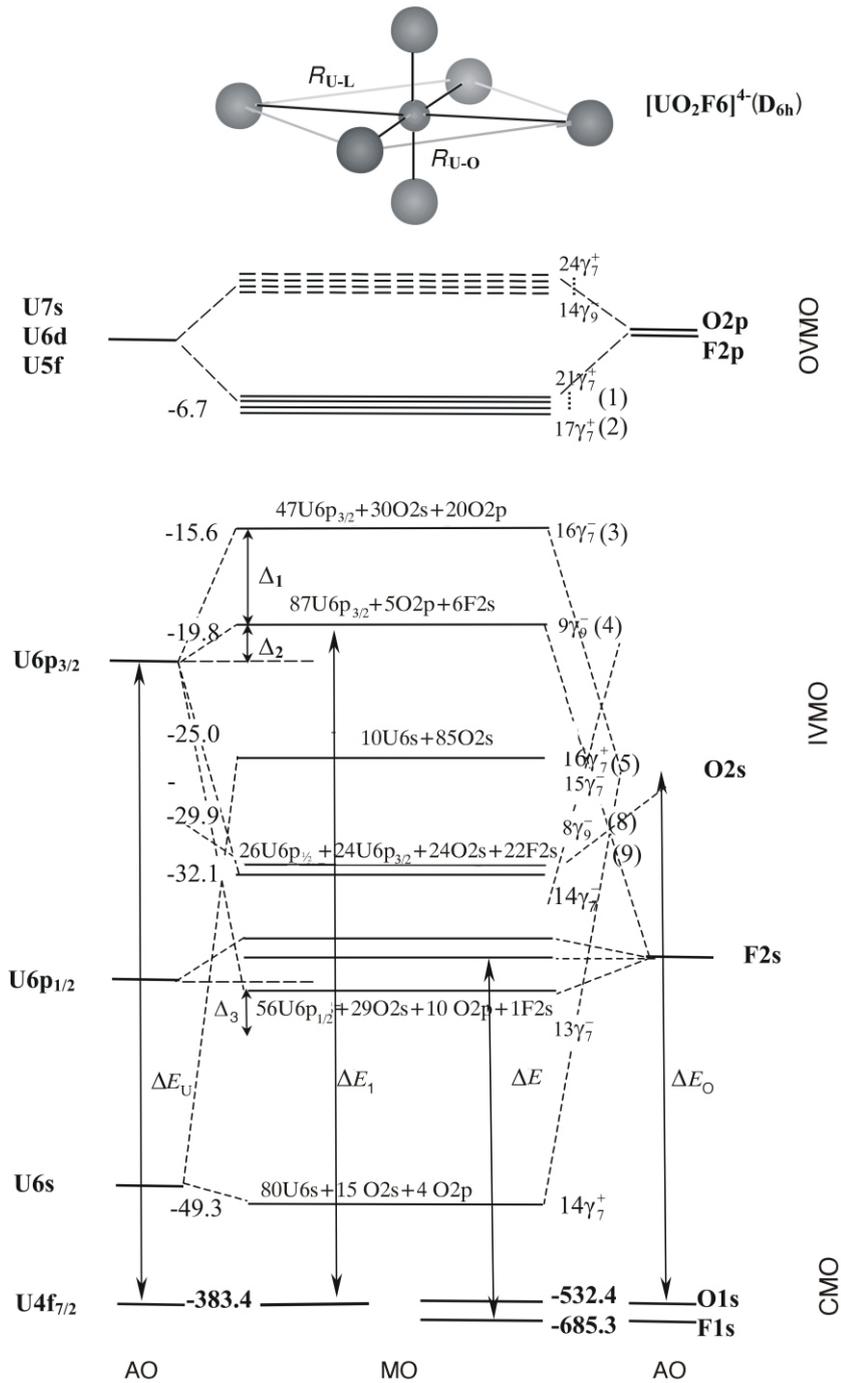


Figure 2. MO schematic diagram for the $[(UO_2)F_6]^{4+}(D_{6h})$ cluster built on the basis of the theoretical and experimental data. Chemical shift is not shown. Arrows indicate some experimentally measurable binding energy differences. Experimental binding energies in eV are given to the right. Energy scale is not kept.

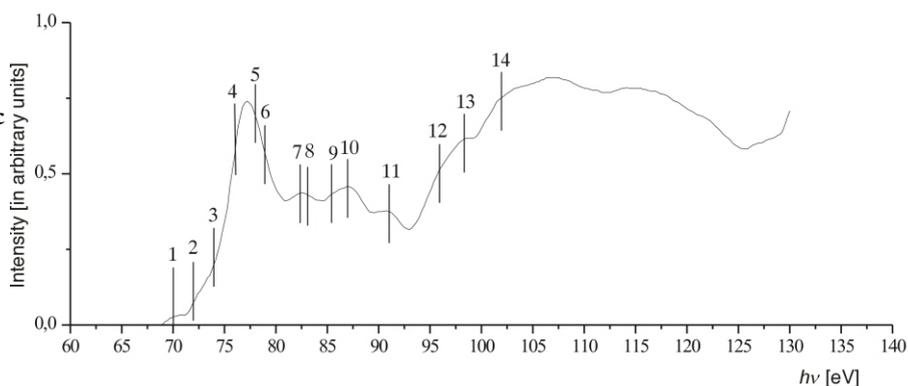
pate in the OVMO formation, while in the IVMO formation these electrons do not participate.

Since quantum mechanics can predict qualitatively the alterations of considered spectral intensities depending on the binding energy, one can vary the initial data for the initial clusters for a better agreement between the theoretical and experimental data. As it has been shown [6, 11], the wider an MO is, the more of its electrons participate (bonding, antibonding) in the chemical bond. The narrower peaks are those attributed to antibonding and quasiatomic orbitals. Therefore, the fact that the $15 \tau^-(7)$ IVMO line is relatively narrow can be explained by an admixture of the $U6p_{3/2}$

and $O2s$ AOs. Despite the fact that it is difficult to establish reliably the $15 \tau^-(7)$ and $13 \tau^-(9)$ IVMO line widths, the obtained data allows the conclusion that the antibonding nature of the $15 \tau^-(7)$ IVMO is lower than the bonding nature of the $13 \tau^-(9)$ IVMO. In other words, electrons from this IVMO couple must bring a positive contribution to the covalent component of the chemical bond in UO_2F_2 .

The most ambiguous for interpretation is the $15 \tau^-(7)$ - $8 \tau^-(8)$ IVMO XPS region. Comparison of the XPS data for UO_2F_2 and UF_4 [6] shows that the quasiatomic $F2s$ AO related peak has to be located at 29.9 eV. Apparently, the calculation was not able to give

Figure 3. X-ray $\text{O}_{4,5}(\text{U})$ emission spectra from UO_2F_2 measured at 3 kV (mA) reflecting the $\text{U}5d \rightarrow \text{U}6p, np, 5f$ electronic transitions. Peak numbers correspond to the transition numbers on the diagram (tab. 3, fig. 4).



the correct binding energies for these MOs. With this in mind this work carried out the decomposition of this spectral region. The knowledge of the correct IVMO sequence is critical for the understanding of the IVMO contributions to the covalent component of the chemical bond in UO_2F_2 . The observed peaks (10) in the spectra can be partially attributed to the shake up process during the photoemission.

One of the experimental confirmations for the IVMO formation in UO_2F_2 is the high resolution X-ray $\text{O}_{4,5}(\text{U})$ emission spectral structure. This spectrum reflecting the $\text{U}5d_{5/2,3/2} \rightarrow \text{U}6p_{3/2,1/2}, np, 5f$ [$\text{O}_{4,5}(\text{U}) \rightarrow \text{P}_{2,3}, \text{O}_6(\text{U})$] electronic transitions was observed in the photon energy range $60 < h\nu < 130$ eV and in the higher energy part interfered with the absorption spectrum, which distorted its structure (fig. 3). During the interpretation of the low resolution X-ray $\text{O}_{4,5}(\text{U})$ emission spectral structure [19] it has been suggested that the low energy structure is associated with the ternary electronic transitions involving the core uranium shells $\text{U}5d_{5/2,3/2} \rightarrow \text{U}6p_{3/2,1/2}$. However, fig. 3 shows that ignoring the participation of the $\text{U}6p_{3/2,1/2}$ electrons in the IVMO formation one can not identify the X-ray $\text{O}_{4,5}(\text{U})$ emission fine spectral

structure. Therefore, the fact that the $\text{U}6p$ shell is not core and participates effectively (experimentally observed) in the IVMO formation was taken into account for a correct interpretation of this structure [11].

The low energy band in the $\text{O}_{4,5}(\text{U})$ XES from UO_2F_2 is wide and poorly resolved despite a relatively high device resolution (fig. 3). As well as for thorium and uranium oxides [20] it can be partially explained by an extra electronic transitions not indicated in the diagram (fig. 4) reflecting the $\text{U}6p_{1/2}, 6s \rightarrow \text{O}(\text{F})2p, 2s$ related IVMOs.

The considered $\text{O}_{4,5}(\text{U})$ XES from UO_2F_2 in the photon energy range 65-110 eV exhibits 14 (not 6 as expected for atomic U) explicit peaks that can be identified using the XPS data and the $\text{U}6p$ and $\text{O}(\text{F})2s$ related OVMO and IVMO concept (tab. 3). The peaks reflecting the transitions from the OVMO and 16 γ^- and 9 γ^- IVMO can be most reliably (practically unambiguously) identified since these peaks are narrow in both XPS and XES. These peaks are: $5d_{5/2} \rightarrow 16 \gamma^-$ (9), $5d_{3/2} \rightarrow 16 \gamma^-$ (11), $5d_{5/2} \rightarrow 9 \gamma^-$ (6), $5d_{3/2} \rightarrow 9 \gamma^-$ (10), $5d_{5/2} \rightarrow \text{OVMO}$ (12), and $5d_{3/2} \rightarrow \text{OVMO}$ (14) (tab. 3). Identification of the peaks reflecting the transitions from the bonding 14 γ^- and 8 γ^- IVMOs is more com-

Figure 4. Schematic diagram of the energy levels for atomic uranium (a) and the $[(\text{UO}_2)\text{F}_6]^{4-}$ (D_{6h}) cluster reflecting uranium close environment in UO_2F_2 (b). Numbers (1) through (11) show the transitions from the IVMO to the core levels. The transitions: $5d_{5/2} \rightarrow \text{OVMO}$ (12), $5d_{5/2} \rightarrow 5f_{7/2}$ (13) and $5d_{3/2} \rightarrow \text{OVMO}$ (14) are not shown (see tab. 3)

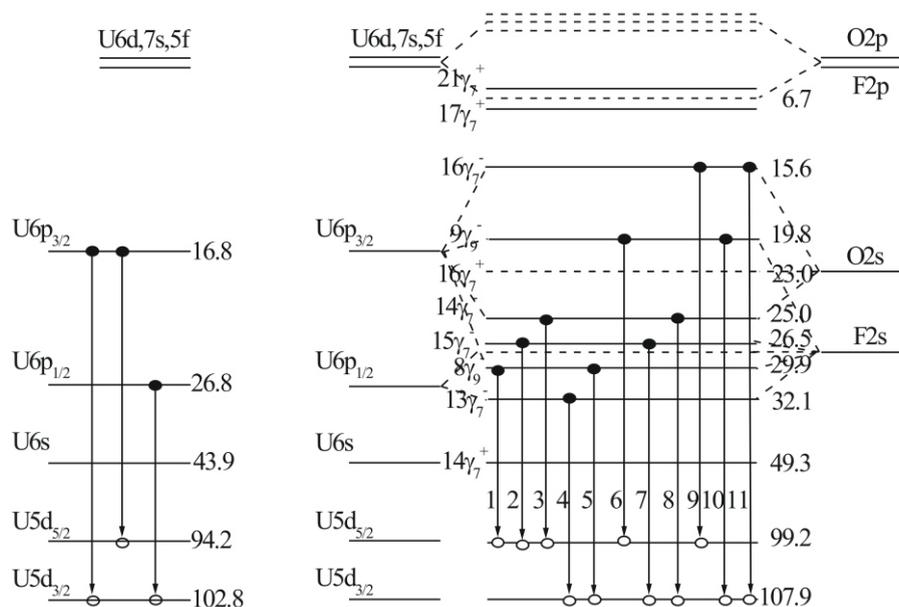


Table 3. X-ray $O_{4,5}(U)$ emission ($U5d \rightarrow U6p, np, 5f$) transition in U and UO_2F_2 on the basis of the XPS and XES data.

U			UO_2F_2		
Transitions	Theory	XPS	Transitions ^(a)	XPS	XES
			(1) $5d_{5/2} \rightarrow 8p^-$	69.3	70.0
			(2) $5d_{5/2} \rightarrow 15p^-$	72.7	72.0
$5d_{3/2} \rightarrow 6p_{1/2}$	78.02	76.0	(3) $5d_{5/2} \rightarrow 14p^-$	74.2	74.0
			(4) $5d_{3/2} \rightarrow 13p^-$	75.8	76.0
$5d_{3/2} \rightarrow 6p_{3/2}$	87.28	86.0	(5) $5d_{3/2} \rightarrow 8p^-$	78.0	78.0
			(6) $5d_{5/2} \rightarrow 9p^-$	79.4	79.0
			(7) $5d_{3/2} \rightarrow 15p^-$	81.4	82.3
$5d_{5/2} \rightarrow 6p_{3/2}$	79.11	77.4	(8) $5d_{3/2} \rightarrow 14p^-$	82.9	83.0
$5d_{5/2} \rightarrow 5f_{5/2}$	97.90	95.4	(9) $5d_{5/2} \rightarrow 16p^-$	83.6	85.4
			(10) $5d_{3/2} \rightarrow 9p^-$	88.1	87.0
$5d_{5/2} \rightarrow 5f_{7/2}$	98.64	104.0	(11) $5d_{3/2} \rightarrow 16p^-$	92.3	91.0
			(12) $5d_{5/2} \rightarrow$ OVMO	92.5	96.0
			(13) $5d_{5/2} \rightarrow 5f_{7/2}$		98.3
$5d_{3/2} \rightarrow 5f_{5/2}$	106.07	104.0	(14) $5d_{3/2} \rightarrow$ OVMO	101.2	102.0

^(a) X-ray $O_{4,5}(U)$ emission peak maxima numbers are given in parentheses (fig. 3)

plicated since the peaks are not that narrow (fig. 3). Despite this some qualitative conclusions can be made. Thus, even the qualitative consideration shows that the low energy region (65–78 eV) of the $O_{4,5}(U)$ XES exhibits the structure, which has to be absent in the case of transitions from only atomic $U6p_{3/2,1/2}$ shells. The peaks in this region can be attributed only to the transitions from the IVMOs of the valence band bottom. Thus, the fine $O_{4,5}(U)$ XES structure is due to the outer and inner valence molecular orbitals built mostly from the $U6p$ and $O(F)2s$ atomic shells.

We would like to note that despite the approximation imperfections, the calculation results for the $[(UO_2)F_6]^{4-}(D_{6h})$ cluster reflecting uranium close environment in UO_2F_2 are in good qualitative agreement with the experimental data, which allowed for the first time a reliable interpretation of the peaks of at least the IVMO ceiling and bottom. These results also can be used for determination of the structure and interatomic distances in uranyl compounds as well as for the interpretation of other X-ray spectra (Auger-, conversion electron, emission, absorption) of uranium compounds.

CONCLUSIONS

Low binding energy X-ray photoelectron and X-ray $O_{4,5}(U)$ emission spectra from UO_2F_2 were measured and interpreted in the relativistic X_α Discrete Variation approximation for the $[(UO_2)F_6]^{4-}(D_{6h})$ cluster reflecting uranium close environment in UO_2F_2 . It yielded a satisfactory qualitative and in some cases quantitative agreement between the theoretical and experimental data.

Despite the traditional opinion that before the chemical bond formation the $An5f$ electrons get promoted to, for example, the $An6d$ atomic orbitals, the theoretical calculations show and experimental data confirm that the $An5f$ atomic shells (more than $U5f$ electrons) can participate directly in the chemical bond formation in uranyl fluorite. The filled $U5f$ states are localized in the OVMO binding energy range from -4 to -9 eV, and the vacant $U5f$ electronic states are generally localized in the low positive energy range (0–6 eV).

The $U6p$ -electrons (about 1 $U6p$ -electron) were experimentally shown to participate significantly in the filled OVMO formation beside the IVMO formation.

The relativistic calculations confirm the fact that the IVMO system associated mostly with the $U6p$ and $O(F)2s$ AOs of uranium and ligands can be subdivided into the two MO groups, one of which ($16p^-$, $14p^-$, $13p^-$) characterizes the bond in the axial direction, and the other one ($9p^-$, $8p^-$) in the equatorial plane. It is important in connection with the development of the experimental technique of determination of the interatomic distances U-L in the axial direction and the equatorial planes of uranyl compounds on the basis of the IVMO XPS parameters.

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**Игор О. УТКИН, Јуриј А. ТЕТЕРИН, Владимир А. ТЕРЕХОВ,
Михаил В. РИЖКОВ, Антон Ј. ТЕТЕРИН, Лабуд ВУКЧЕВИЋ**

**ПРОУЧАВАЊЕ ЕЛЕКТРОНСКЕ СТРУКТУРЕ УРАНИЈУМФЛУОРИТА UO₂F₂
РЕНДГЕНСКОМ ФОТОЕЛЕКТРОНСКОМ СПЕКТРОСКОПИЈОМ**

У раду се разматра фина структура рендгенског фотоелектронског спектра слабо везаних електрона (0-40 eV) и структура рендгенског O_{4,5}(U) емисионог спектра из UO₂F₂. Коришћен је релативистички X_α дискретни варијациони поступак за прорачун [(UO₂)F₆]⁴⁻ (D_{6h}) кластера који одражава непосредну околину уранијума у UO₂F₂. Показано је да U5f електрони непосредно учествују у образовању хемијске везе. Такође је показано да U6p електрони учествују у образовању не само унутрашњих, већ и спољашњих валентних молекулских орбитала. Утврђен је секвенцијални поредак унутрашњих валентних молекулских орбитала у подручју енергија везе од 12-40 eV. Ово је од значаја за развој поступка одређивања међуатомских растојања уранијумских једињења у аксијалном правцу и екваторијалној равни, заснованог на рендгенској фотоелектронској спектроскопији.